

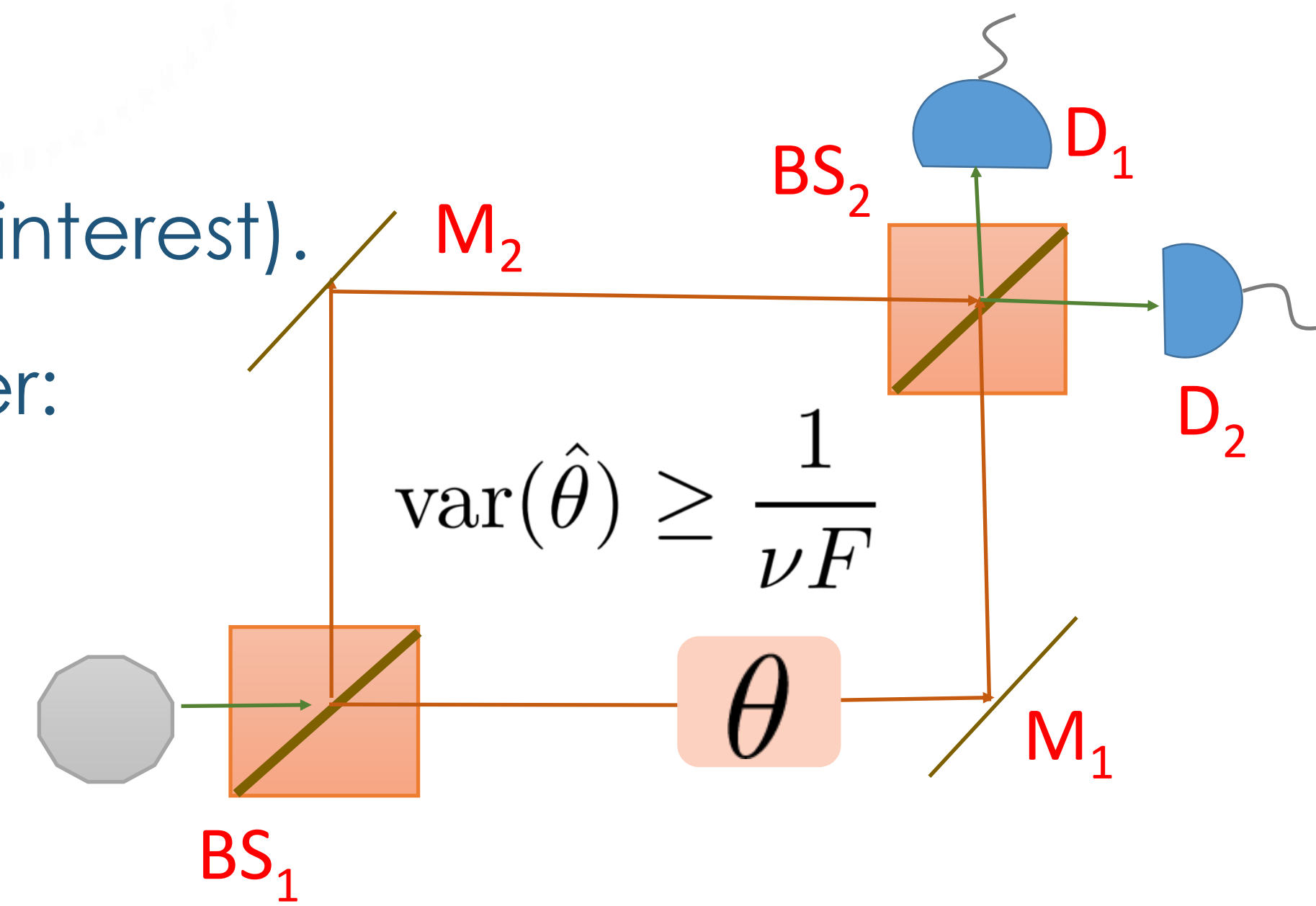
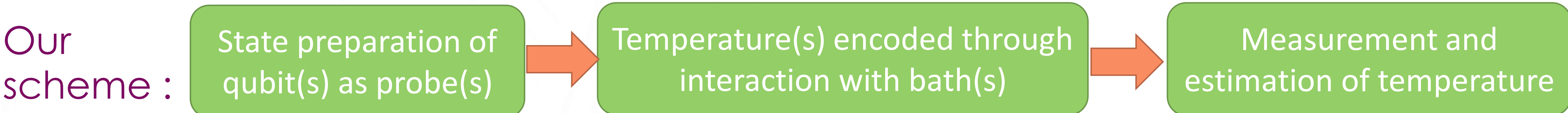
INTRODUCTION



- Question:** What does it mean to have a superposition of thermalizing quantum channels?
- Approach:**
 - Exert Q-ctrl over the state of a bath – prepare it in a superposition of purified Gibbs' state.
 - Exert Q-ctrl over the path of the probe, prepare it in a superposition where it goes through 2 different baths.
- Refined question:** How effectively can we measure the temperature of the bath(s) in the above 2 approaches?

MULTIPARAMETER ESTIMATION & CR BOUND(S)

- How well can a parameter be estimated?
- Density matrix – contains all the information we need (nonlinear fn. of parameter of interest).
- QFI: Mostly used in context of phase estimation – for e.g. in case of MZ interferometer:



$$\text{var}(\hat{\theta}) \geq \frac{1}{\nu F}$$

THERMALIZATION MODEL

Purified thermal bath constituted by 2 qubits:

$$|\theta^\beta\rangle \equiv \frac{e^{-E_0\beta/2}}{\sqrt{e^{-E_0\beta} + e^{-E_1\beta}}} |0,0\rangle + \frac{e^{-E_1\beta/2}}{\sqrt{e^{-E_0\beta} + e^{-E_1\beta}}} |1,1\rangle$$

Interaction unitary (U):

$$U_{PB}^\eta = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \sqrt{1-\eta} & \sqrt{\eta} & 0 \\ 0 & -\sqrt{\eta} & \sqrt{1-\eta} & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$T = 1/\beta$

$$H_B = H_0 \otimes \mathbb{I}_2 + \mathbb{I}_2 \otimes H_0$$

$$H_0 = E_0|0\rangle\langle 0| + E_1|1\rangle\langle 1|$$

ν – Repetitions of expt.
 F – Q. Fisher Information

Multiparameter CR bound:

$$\text{Cov}(\vec{T}) \geq \frac{Q_T^{-1}}{\nu}$$

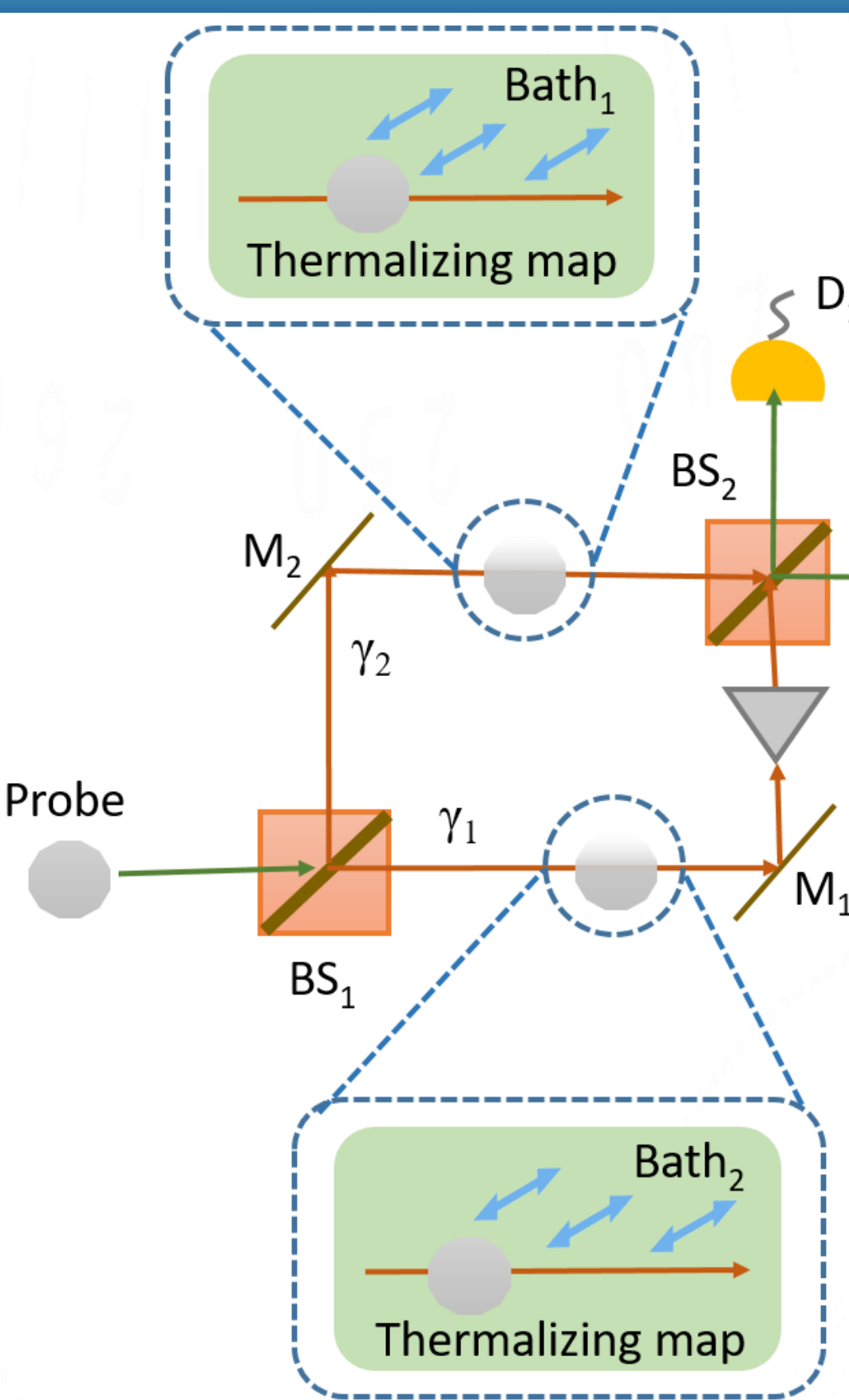
Bounds on variances obtained from the above matrix inequality:

$$\text{Var}(T_1) \geq \frac{Q_{T_2 T_2}}{\nu \det(Q_T)}$$

$$\text{Var}(T_2) \geq \frac{Q_{T_1 T_1}}{\nu \det(Q_T)}$$

$$\left(\text{Var}(T_1) - \frac{Q_{T_2 T_2}}{\nu \det(Q_T)} \right) \left(\text{Var}(T_2) - \frac{Q_{T_1 T_1}}{\nu \det(Q_T)} \right) \geq \left[\text{Cov}(T_1, T_2) + \frac{Q_{T_1 T_2}}{\nu \det(Q_T)} \right]^2$$

NUMERICAL RESULTS FOR THE MACH ZEHNDER SETUP

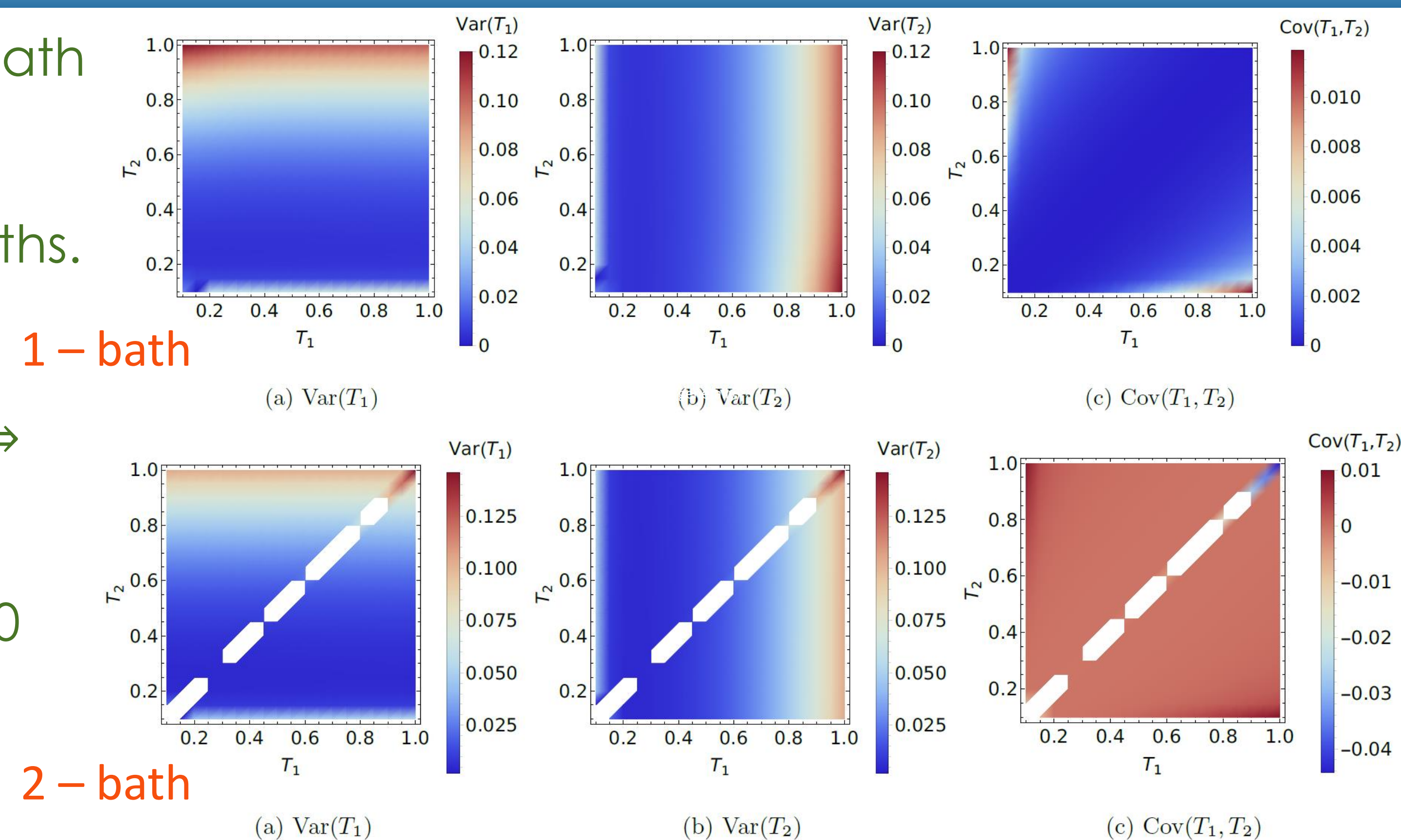


- 1-bath:** The state of the (thermal) bath depends on the control (path).
- 2-bath:** Two independent baths on the paths.

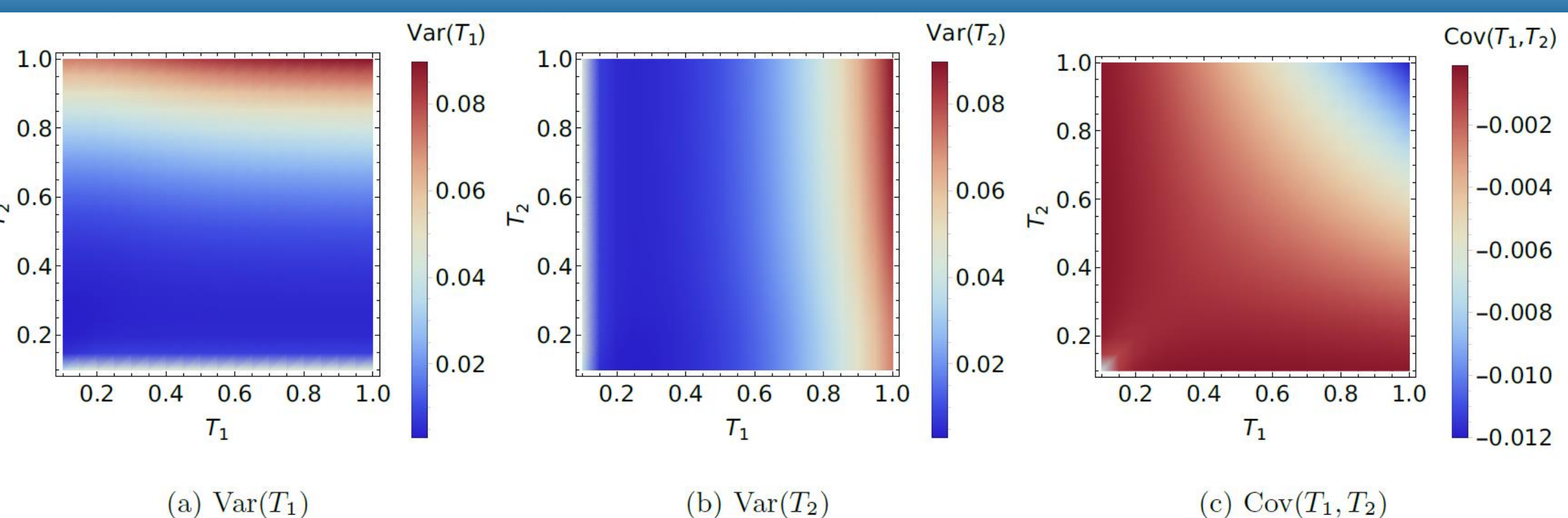
For a single qubit as a probe:

- 1-bath/ 2-bath : $|QFIM| = 0 \Rightarrow$ Bounds cannot be obtained.
- For a two-qubit probe, $|QFIM| \neq 0$ (generally).

$$\rho_{\text{in}} = |0\rangle\langle 0| \otimes |0\rangle\langle 0| \quad \Phi = \pi/2, \nu = 100$$



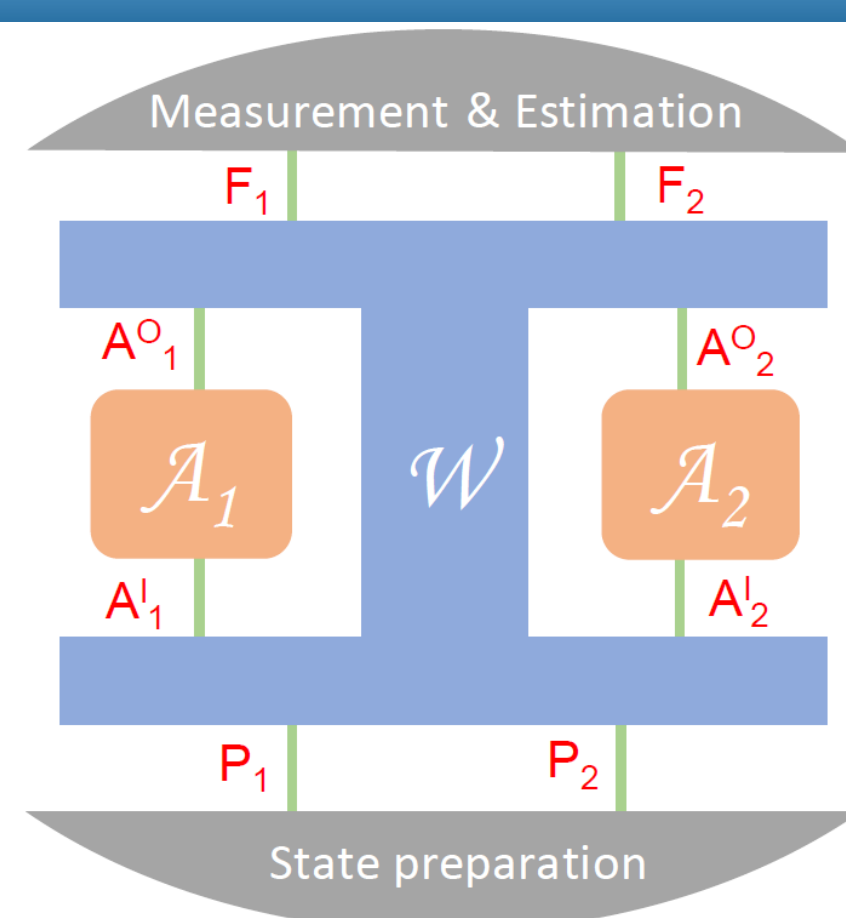
NUMERICAL RESULTS FOR THE QUANTUM SWITCH BASED SETUP



- Indefinite causal order via. Q - Ctrl over order of application of 2 thermalizing channels.

$$\rho_{\text{in}} = |0\rangle\langle 0| \otimes |+\rangle\langle +| \quad \nu = 100$$

$\mathcal{A}_1, \mathcal{A}_2$: Quantum channels thermalizing to distinct temperatures.



CONCLUSIONS

- Mach-Zehnder interferometer can be used for the task of estimating 2 temperatures, with 2 qubits as a probe.
- Qualitatively similar results for setups based on other quantum processes.
- Parallel scheme with simultaneous measurement of two temperatures using two independent probes outperforms all the others.

REFERENCES

[1] Araújo, M., Feix, A., Navascués, M., & Brukner, C. A purification postulate for quantum mechanics with indefinite causal order. *Quantum*, 1, 1–13 (2017).
 [2] Oi, D. K. L. Interference of Quantum channels. *Phys. Rev. Lett.*, 91(6), 67902 (2003).
 [3] Liu, J., Yuan, H., Lu, X.-M., & Wang, X. Quantum Fisher Information Matrix and Multiparameter Estimation. *Journal of Physics A: Mathematical and Theoretical*, 53(2), 23001 (2019).
 [4] Mukhopadhyay, C., Gupta, M. K., & Pati, A. K. Superposition of causal order as a metrological resource for quantum thermometry. *arXiv:1812.07508v1* (2018).

